

QGS : Quantum groups seminar

Quantum affine algebras and spectral k -matrices

1. The Yang-Baxter equation (YBE)

$$R(z) \in \mathcal{A} \otimes \mathcal{A}((z)) \text{ s.t.} \quad \text{some algebra}$$

$$\begin{aligned} R_{12}(z) R_{13}(zw) R_{23}(w) &= \\ &= R_{23}(w) R_{13}(zw) R_{12}(z) \end{aligned}$$

$$R_{12}(z) = R(z) \otimes 1 \quad R_{23}(z) = 1 \otimes R(z)$$

$$R_{13}(z) = (\text{flip} \otimes 1)(R_{23}(z))$$

$$R(z) = \text{R-matrix with spectral parameter } z$$

2. Quantum integrable systems

$$\mathcal{A} = \text{end}(\mathbb{C}^2)$$

$q \in \mathbb{C}^\times$ - quantum parameter

$$R(z) := \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{q^{-1}(z-1)}{z-q^{-2}} & \frac{1-q^{-2}}{z-q^{-2}} & 0 \\ 0 & \frac{z(1-q^{-2})}{z-q^{-2}} & \frac{q^{-1}(z-1)}{z-q^{-2}} & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

XXZ $\frac{1}{2}$ -spin
Heisenberg model / 6-vertex
model

$$\rightsquigarrow U_q \widehat{\mathfrak{sl}}_2$$



$$\mathbb{C}^2$$

3. Quantum affine algebras

$\mathfrak{g}$ = simple f.d. Lie algebra / $\mathbb{C}$

$$\widehat{\mathfrak{g}} = \underbrace{\mathfrak{g} \otimes \mathbb{C}[t, t^{-1}]}_{\mathcal{L}\mathfrak{g}} \oplus \mathbb{C}^c \oplus \mathbb{C}^d$$

↗ affine Lie algebra
↖ loop algebra
↖ central
↗ derivation

(Kac-Moody algebra)

$$\widehat{\mathfrak{sl}}_2$$


$$\begin{aligned} e_1, f_1 &\rightsquigarrow e \otimes 1, f \otimes 1 \\ e_0, f_0 &\rightsquigarrow f \otimes t, e \otimes t^{-1} \end{aligned}$$

$$\widehat{U_q \mathfrak{g}} \supseteq (U_q \widehat{\mathfrak{g}})' \xrightarrow{q^c=1} (U_q \widehat{\mathfrak{g}})' \simeq U_q \mathcal{L}\mathfrak{g}$$

↗ Drinfeld-Jimbo quantum group

↖ Quantum loop algebra

(usual Kac-Moody presentation)

(Drinfeld loop presentation)

4. Representations

$$U_q \mathcal{O}$$



$$\mathcal{O}^{\text{int}}$$

- $\dim(V) = \infty$
- semisimple
- $\underline{\lambda} \in \mathbb{N}^I$
- braided

$$U_q \mathcal{L} \mathcal{O}$$



$$\text{Rep}^{\text{fd}}(U_q \mathcal{L} \mathcal{O})$$

- $\dim(V) < \infty$
- non-semisimple!
- $P(u) \in \mathbb{C}[u]^I$
- non-braided!

Remark $U_q \mathcal{L} \mathcal{O}$ is $\mathbb{Z}$ -graded ($\deg(e_i/f_i) = 0$, $\deg(e_0/f_0) = \pm 1$) and endowed with

$$(a \in \mathbb{C}^\times) \quad \tau_a \hookrightarrow U_q \mathcal{L} \mathcal{O}$$

$$\tau_a(x) = a^m x \quad \text{if } \deg(x) = m$$

$$V \in \text{Rep}^{\text{fd}}(U_q \mathcal{L} \mathcal{O}) \rightsquigarrow V(a) := \tau_a^*(V)$$

Thm (Drinfeld)

$$\exists R(z) \in (U_q \mathcal{L} \mathfrak{g} \hat{\otimes} U_q \mathcal{L} \mathfrak{g})[[z]] \text{ s.t.}$$

- $R_{VW}(z) \in \text{end}(V \otimes W)[[z]]$
- $\text{flip} \circ R_{VW}(z)$ **intertwiner**
- $\Delta \otimes 1(R(z)) = R_{13}(z) R_{23}(z)$
 $1 \otimes \Delta(R(z)) = R_{13}(z) R_{12}(z)$

$\leadsto R_{VW}(z)$ satisfies (YBE).

idea: $U_q \hat{\mathfrak{g}}$ is quasi-triangular Hopf algebra

$$R = q^{\Omega_0} \cdot \sum_{\mu > 0} e_{\mu} \otimes f_{\mu} \in U_q \hat{\mathfrak{n}}_+ \hat{\otimes} U_q \hat{\mathfrak{n}}_-$$

$$\begin{aligned} \Omega_0 &= \sum u_i \otimes u^i \in \mathfrak{h} \\ &\quad + c \otimes d + d \otimes c \end{aligned}$$

$$\leadsto R_{VW}(z) := \tau_z \otimes 1(R) \in V \otimes W$$

5. Reflection equation

From integrable models with boundary conditions (Sklyanin, Cherednik, ...)

$k(z) \in \mathcal{A}((z))$ s.t.

$$\begin{aligned} k_1(z) R_{21}^{\varphi}(zw) k_2(w) R(w/z) \\ = R_{21}^{\varphi\varphi}(w/z) k_2(z) R^{\varphi}(zw) k_1(w) \end{aligned}$$

$\leadsto k(z) =$ k -matrix with
spectral parameter z

examples

- many computational examples for $U_q \mathfrak{Lg}$
(e.g. Delius - George, Regelski - Vlaar ...)
- universal **constant** solutions for
finite-type $U_q \mathfrak{g}$ (Bao - Wang,
Balagovic - Kolb)

6. Quantum symmetric pairs

In both cases, the construction of the k -matrix is deeply related to Letzter-Kolb quantum symmetric pairs

$$\mathcal{B} \subseteq U_q \mathfrak{g} \leftarrow \mathfrak{g} \text{ Kac-Moody}$$

$$\swarrow \text{coideal } \Delta(\mathcal{B}) \subseteq \mathcal{B} \otimes U_q \mathfrak{g}$$

$$(q=1) \quad \mathcal{B} \longrightarrow U(\mathfrak{g}^\theta)$$

$$\mathcal{B} = \text{Ad}(w_X) \circ \omega \circ \tau$$

finite-type Dynkin
subdiagram

diagram automorphism
with $\tau^2 = 1$ and

$\tau|_X =$ action induced
by w_X

longest element in W_X

$B \subseteq U_q \mathfrak{g}$ controls the intertwining property of k on $V \in \mathcal{O}^{int}$:

$$\forall v \in V, b \in B$$

$$k_v \cdot b \cdot v = \varphi(b) \cdot k_v \cdot v$$

↑ some diagram automorphism

$$\Rightarrow k_v : V \xrightarrow{\sim} V^\varphi \text{ as } B\text{-modules}$$

In affine type, the intertwining equation involves the *parameter inversion*, i.e.

$$k_v(z) : V(z) \longrightarrow V(1/z)^\varphi$$

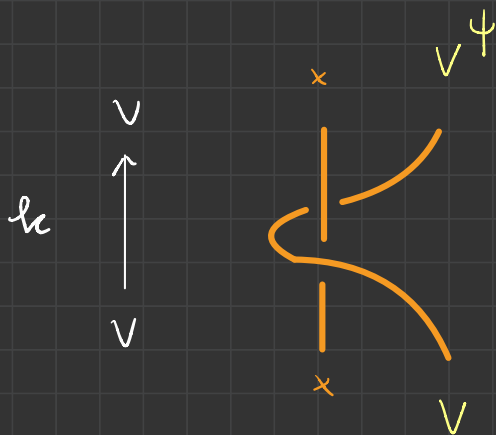
Goal: obtain *spectral k -matrices* from constant (Kac-Moody) k -matrices.

Big problem: φ cannot perform the parameter inversion

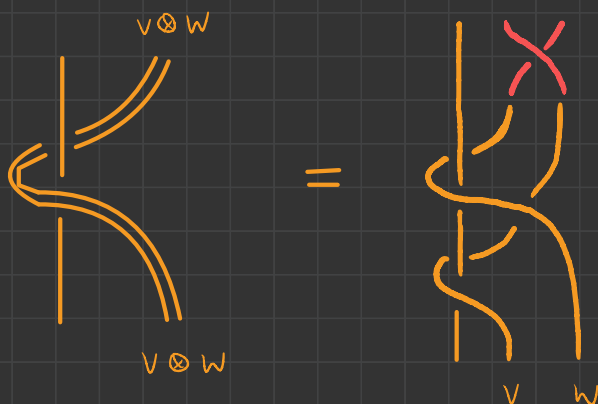
7. A new approach (A. - Vlaar)

- $(H, R) = \text{qt Hopf algebra}$
- $\psi: H \rightarrow H = \text{algebra automorphism}$
- $V \in \text{Rep}(H) =: \mathcal{C}$

We would like an operator $k: V \rightarrow V^\psi$



and we want to
get the RE from
a suitable coupling
identity



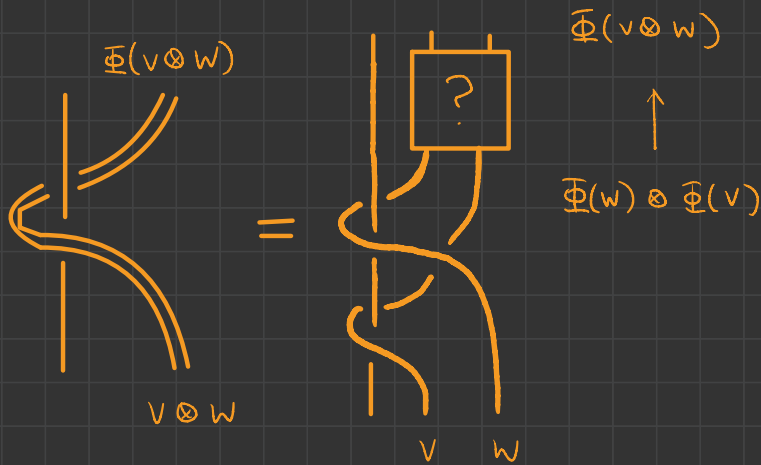
← not very
canonical

$$\psi = \text{id}, \varphi$$

Assume instead we have a functor

$\Phi: \mathcal{C} \rightarrow \mathcal{C}$ codifying the twisting so that

$k: V \rightarrow \Phi(V)$ (say $\Phi = \varphi^*$) and



Our choice becomes almost canonical once we

assume Φ is a braided functor

$$(\Phi, \gamma): \mathcal{C} \rightarrow \mathcal{C}^{\text{op}}$$

← opposite braided
monoidal category

$$\gamma: \Phi(V \otimes W) \xrightarrow{\sim} \Phi(W) \otimes \Phi(V)$$

← precisely what we need !

Algebraically, our new setting is given by the following data:

- (H, R) = quasitriangular Hopf algebra
- $\psi: H \rightarrow H$ = algebra automorphism
- $J \in H \otimes H$ = Drinfeld twist

satisfying

$$(H^{\psi})^{\psi} = H_J \quad R_{21}^{\psi\psi} = R_J$$

where

$(H^{\psi}, R^{\psi\psi})$ is the ψ -twisted qtHA

$$\Delta^{\psi}(x) := \psi \otimes \psi \circ \Delta \circ \psi^{-1}(x)$$

$$R^{\psi\psi} = \psi \otimes \psi(R)$$

(H_J, R_J) is the J -twisted qtHA

$$\Delta_J(x) := J \cdot \Delta(x) \cdot J^{-1}$$

$$R_J := J_{21} \cdot R \cdot J^{-1}$$

We assume (H, R) is fixed and we call (ψ, J) a **twist pair**. Note that this is equivalent to have a **braided equivalence**

$$(\psi^*, J) : \text{Rep}(H, R) \longrightarrow \text{Rep}(H, R)^{\text{op}} \\ \parallel \\ \text{Rep}(H^{\text{cop}}, R_{21})$$

Finally we get the **k-matrix**.

Def A **cylindrical bialgebra** is a tuple (H, R, ψ, J, k) as before with $k \in H$ s.t.

$$\Delta(k) = J^{-1} \cdot 1 \otimes k \cdot R^+ \cdot k \otimes 1$$

↖
new coupling identity

Cylindrical bialgebras provide the right framework to study a more general reflection equation:

$$\begin{aligned}
 (1) \quad R_{21}^{\psi\psi} \cdot 1 \otimes k \cdot R^{\psi} \cdot k \otimes 1 &= \\
 &= J_{21} \cdot R \cdot \underbrace{\bar{J}^{-1} \cdot 1 \otimes k \cdot R^{\psi} \cdot k \otimes 1}_{\Delta(k)} = \\
 &= k \otimes 1 \cdot R_{21}^{\psi} \cdot 1 \otimes k \cdot R
 \end{aligned}$$

(2) k, R induce an external action of the cylindrical braid group on $V^{\otimes N}$

(3) Set $\vartheta(x) := k^{-1} \cdot \psi(x) \cdot k$ and let $B_k \subseteq H^{\vartheta}$ be the maximal coideal subalgebra. Then, $\text{Rep}(B_k)$ is a braided module category over $\text{Rep}(H)$.

8. New k -matrices for $U_q \mathfrak{g}$

Let $U_q \mathfrak{g}$ be a quantum Kac-Moody algebra and $B \subseteq U_q \mathfrak{g}$ a QSP corresponding to $\vartheta = \text{Ad}(w_x) \circ w \circ \tau$.

The construction of a k -matrix for $U_q \mathfrak{g}$ supported on B requires the use of quantum Weyl group operators:

$T_i = q$ -analogue of triple exponentials

$[L, KR, LS]$ action of B_W on $U_q \mathfrak{g}$
and $\mathcal{O}^{\text{int}}$ representations

[Kamnitzer-Tingley] $\exists t_x := (\text{Cartan}) T_{w_x}$

s.t.

$$R_x := t_x^{-1} \otimes t_x^{-1} \Delta(t_x)$$

t_x is a half-balance for $U_q \mathfrak{g}_x$

Modified quantum Weyl group operators

Given **Satake** diagrams (X, τ) and (Y, η)
define suitable corrections

$$T_{X, \tau} := (\text{Cartan}) \cdot T_{W_X}$$

$$R_{X, \tau} := (\text{Cartan}) \cdot R_X$$

and we set

$$\vartheta_q(X, \tau) := \text{Ad}(T_{X, \tau}) \circ \omega \circ \tau \hookrightarrow U_{qg}$$

$$\psi_{Y, \eta} := \text{Ad}(T_{Y, \eta}) \circ \omega \circ \tau$$

such that $(\psi_{Y, \eta}, R_{Y, \eta})$ is a twist pair.

Rem in finite type, we choose $Y = I$ and
 $\psi = \text{op}_I \circ \tau$ (cf. Balagovic-Kolb).

Finally we can state our

Thm (A. - Vlaar)

There exists a canonical series $\mathcal{X}_{X,\tau} = 1 + \sum_{\mu} \mathcal{X}_{\mu}$
with $\mathcal{X}_{\mu} \in (U_q \mathfrak{n}_+)_{\mu}$ s.t.

$$k_{Y,\eta} := \begin{pmatrix} T_{Y,\eta}^{-1} & T_{X,\tau} \end{pmatrix} \circ (\text{Cartan}) \circ \mathcal{X}_{X,\tau}$$

is a k -matrix for the twist pair $(\psi_{Y,\eta}, R_{Y,\eta})$
with support $B_{X,\tau}$.

Affine type $Y = I \setminus \{0, \tau(0)\}$ and $\eta(0) = \tau(0)$

$(0 \notin X)$ and $\eta|_Y = \text{id}_Y$. For any

$V \in \text{Rep}^{\text{fd}} U_q \mathfrak{L}_{\mathfrak{g}}$, we get a formal series

$$k_{Y,\eta}(z) \in \text{end}(V)[[z]]$$

satisfying the spectral reflection equation
in $\psi_{Y,\eta}$.